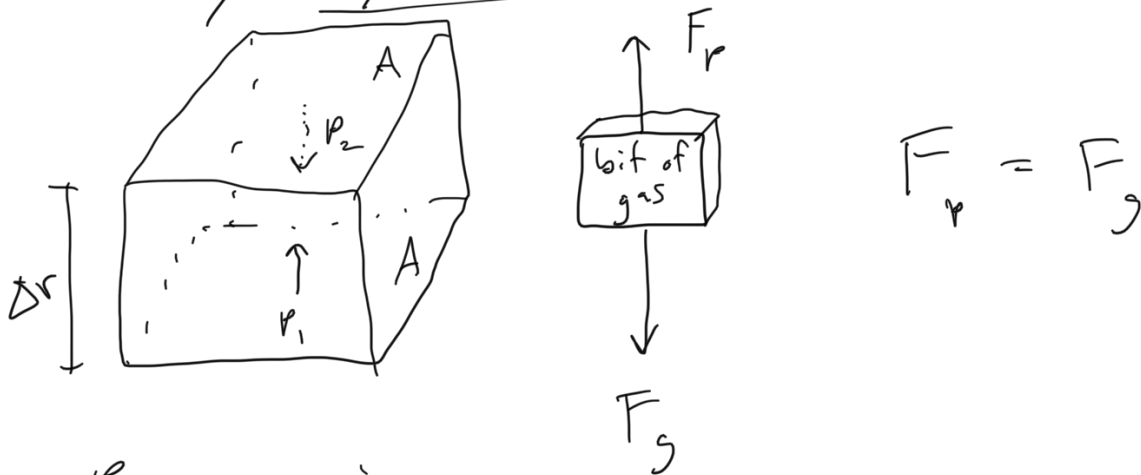


ASTR 2500 - Ch.14 & 15

Stellar Atmospheres & Classification

Atmospheres in general (that of the Earth, Sun, etc.) have a structure determined by hydrostatic equilibrium



Pressure is a
force per area

To keep the gas stationary, the difference in pressure across the box must equal F_g

$$\underbrace{(p_2 - p_1)}_{\text{difference}} A = - \frac{GM(r) \overbrace{\rho A \Delta r}^m}{r^2}$$

If box
small
enough

$$\Delta P = - \frac{GM\rho}{r^2} dr$$

$$\frac{dP}{dr} = - \frac{GM(<r)\rho}{r^2}$$

pressure
gradient

gravity

★ For a spherical dist. w/ density that only depends on radius, all mass interior can be thought of as in the center, and the grav. force of all mass exterior cancel out, hence $M(<r)$

Luminosity Classes

- w/o distance hard to get class, but turns out line widths depend on class

- narrower lines in I (supergiant) than V (dwarf / MS) : why?

Surface gravity, g , is \downarrow bc $R \uparrow$,
1. on the

∴ thus the pressure is lower, ∴
 lines have less pressure broadening

$$\frac{dP}{dr} = - \frac{GM\rho}{r^2} = -g\rho$$

Line comes from a shell dr that corresponds to the optical depth

$$\tau = 0 \rightarrow 1, \quad \tau = n\sigma x, \text{ or here,}$$

can write as $d\tau = -n\sigma dr$ ($r \downarrow$ as $\tau \uparrow$)

Can rewrite the cross-section as a quantity per mass (in this shell), called the

opacity: $K = \frac{\sigma}{\text{mass}}, \text{ or } \rho K = n\sigma$

$$d\tau = -\rho K dr$$

$$\frac{dP}{dr} = \frac{dP}{-d\tau/\rho K} = -g\rho$$

$$\frac{dP}{d\tau} = \frac{g}{K} \rightarrow \rho \sim \frac{g}{K} \tau$$

Since $\tau \sim 1$, $\rho \sim \frac{g}{K}$ ∴ K is rough constant in stars, so $\downarrow g$, $\downarrow \rho$, narrower lines

★ Also, gas in stars generally follows the ideal gas law

$$P = nkT = \frac{\rho kT}{\mu m_p}$$

Density is $\frac{M}{V}$, but particles have diff. individual masses, so use 3 categories

$$\rho = \rho_H + \rho_{He} + \rho_{\text{metal}}$$

→ all elements w/ $Z \geq 2$ are called "metals", even though O, N, Ne, etc. are not technically metals

Express the ratio of each cat. by the total density, define

$$X \equiv \rho_H / \rho$$

$$Y \equiv \rho_{He} / \rho$$

$$Z \equiv \rho_{\text{metal}} / \rho = 1 - X - Y$$

$$\tau = \rho_{\text{metal}} / \rho$$

For the Sun, we observe

$$X_0 = 0.734, \quad Y_0 = 0.250, \quad Z_0 = 0.016$$

This is very similar to the primordial abundance, $X \sim 0.75 / Y \sim 0.25 / Z = 0$

Often we care about # density

$$n = \frac{\rho}{\mu m_p}$$

$$\mu_H = \begin{cases} 1 & \text{neutral} \\ \frac{1}{2} & \text{ionized} \end{cases}, \quad \mu_{He} = \begin{cases} 4 & \text{neutral} \\ \frac{4}{3} & \text{ionized} \end{cases}$$

$$\mu(\text{ionized}) = \frac{\rho}{n m_p}$$

$$= \left(2X + \frac{3}{4}Y + \frac{1}{2}Z \right)^{-1}$$

$$\mu_0(\text{ion.}) = 0.60$$

Stellar Interiors

$\tau = 1.1 \dots 1.1 \dots 1.1 \dots$

Just like atmospheres, the interior structure is governed by HSE

$$\frac{dP}{dr} = - \frac{G M(r) \rho(r)}{r^2}$$

By furthering calculus to our ends,
can rewrite as

$$\frac{\Delta P}{\Delta r} \approx - \frac{G \langle M \rangle \langle \rho \rangle}{\langle r \rangle^2}$$

$$\langle M \rangle \sim \frac{M_0}{2}$$

$$\langle \rho \rangle \sim \frac{M_0}{\frac{4}{3} \pi R_0^3} = \rho_0 \approx 1400 \frac{\text{kg}}{\text{m}^3}$$

$\sim 1.4 \times \rho_{\text{H}_2\text{O}}$

$$\langle r \rangle \sim \frac{R_0}{2}$$

$$\Delta P = P_{\text{surf}} - P_c = 0 - P_c$$

$$\Delta r = R_0 - 0$$



$$P_c \approx \frac{P_{\text{ET}}}{3} G \rho_0^2 R_0^2$$

What about T_c ? Can use

ideal gas law!

$$P(r) = \frac{\rho(r) k T(r)}{\mu m_p}$$

$$T_c = \frac{2 G M_0 \mu_0 m_p}{R_0 k}$$

$$\mu_0 \approx 0.60$$

$$T_c \approx 3 \times 10^7 \text{ K}$$

(Correct analysis gives $2 \rightarrow \downarrow T_c$)

For stars on the main sequence

$$M / M_0 < 1.66 M_0$$

$$R / R_0 = 1.06 (M / M_0)^{0.945}$$

$$\approx M / M_0$$

Since $T_c \propto \frac{M}{R} \mu$, + μ same for any star (since H dominates mass)

$$T_c \propto 0.6 \frac{M}{R} \propto 0.6 \frac{M}{\mu} \frac{M_0}{R_0}$$

$\propto \text{const}$

This const. temperature determines
the main sequence — what sets
★ T_c ?

Energy Generation

- in the 1800s, no nuclear physics (atomic structure unknown)
- gravity understood (Newtonian), so Helmholtz & Kelvin considered the formation of the Sun from a contracting gas cloud — its light would then be the conversion of grav. pot. E to heat

For a spherical object,

$$U = -q \frac{GM^2}{R} \quad (q \text{ depends on } \rho(r))$$

This is the E that must be removed to collapse an object from $r = \infty$ to $r = R$ that has mass M

If the Sun's light comes from this energy, & we assume it has been emitting it at a constant rate...

$$U \rightarrow J \quad L \rightarrow J s^{-1}$$

Can estimate the total lifetime of the Sun!

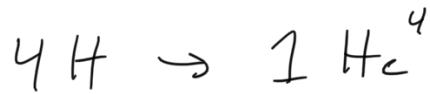
$$t_{\text{KH}} \approx \frac{U_0}{L_0} \approx \underline{50 \text{ Myr}}$$

Geologists thought the Earth was older, but Kelvin used this estimate to cow them, setting back geology as a field for decades - almost 100yr!

Moral: Physicists who don't respect the science of other fields are evil
↳ e.g., physicists making epidemiological models for COVID-19 spread at same university

In the 1930s, nuclear fusion
source

recognized as energy)



$$\begin{aligned} 4m_p &= 6.6905 \times 10^{-27} \text{ kg} \\ 1m_{\text{He}} &= 6.6447 \times 10^{-27} \text{ kg} \end{aligned} \quad \left. \vphantom{\begin{aligned} 4m_p \\ 1m_{\text{He}} \end{aligned}} \right\} 0.0458 \times 10^{-27} \text{ kg} \text{ difference}$$

Where does mass go? $\$$

$$E = mc^2!$$

$$E = \Delta m c^2 = 4.1 \times 10^{-12} \text{ J}$$

Small amount of E , but Sun has lots of protons! Assuming it's entirely H

$$N_{\text{H}} = \frac{M_{\odot}}{m_p} \approx \frac{2 \times 10^{30} \text{ kg}}{1.67 \times 10^{-27} \text{ kg}} \approx 10^{57}$$

If fuse all into He ($4\text{H} \rightarrow 1\text{He}$),

then have Δm per reaction, &

there are $\frac{N_{\text{H}}}{4}$ reactions, releasing

$$E_{\text{fus}} = \frac{N_{\text{H}}}{4} \Delta E = \frac{10^{57}}{4} 4 \times 10^{-12} \text{ J} = \underline{10^{45} \text{ J}}$$

This is $\sim 2000 \times$ more E than provided

$$\text{by } U_0, \text{ so } t_{\text{fus}} \approx \frac{E_{\text{fus}}}{L_0} \approx 100 \text{ Gyr}$$

Stars like the Sun only convert $\sim 10\%$ of their H to He, so actual lifetime is $10^{10} \text{ yr} = \underline{10 \text{ Gyr}}$

$$\text{Since } \tau \propto \frac{M}{L} \text{ \& } L \propto M^4,$$

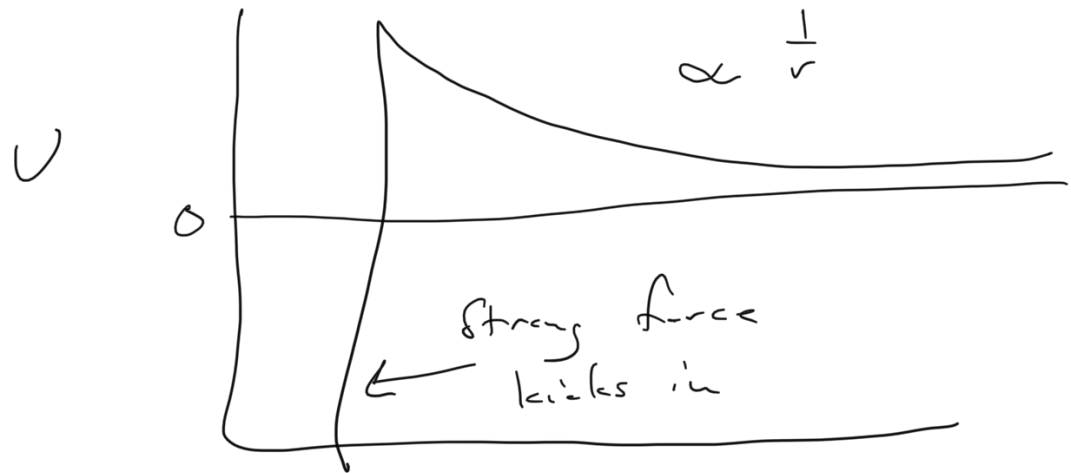
$$\tau \approx 10 \text{ Gyr} \left(\frac{M}{M_{\odot}} \right)^{-3}$$

for main sequence stars

Fusion reactions

How can 2 protons fuse? The repulsive Coulomb force is $\propto \frac{1}{r^2}$ - if protons were point-like, the force would repel them always

But p^+ aren't - they're made of quarks governed by the strong force, which is VERY attractive



Need enough energy to get close enough for strong force to take over
 $r \sim 10^{-15} \text{ m}$

$$U \approx \frac{e^2}{4\pi\epsilon_0 r} \approx 1.4 \text{ MeV}$$

Do protons in the center of the Sun have this much kinetic E ?



$$\langle E \rangle = \frac{3}{2} kT_c \approx 2 \text{ keV}$$

Nope! So how does fusion occur?

Particles are not just particles, but also waves (QM saves the day!)

$$\lambda_{\text{prot}} = \frac{h}{p} \quad , \quad p = m_p v$$

$$\lambda_{\text{prot}} \approx 10^{-13} \text{ m}$$

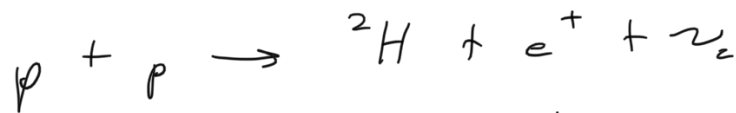
$$\sim 100 r_{\text{SF}}$$

Essentially, the proton's position is uncertain, & has a chance of getting close enough even though classically that's impossible

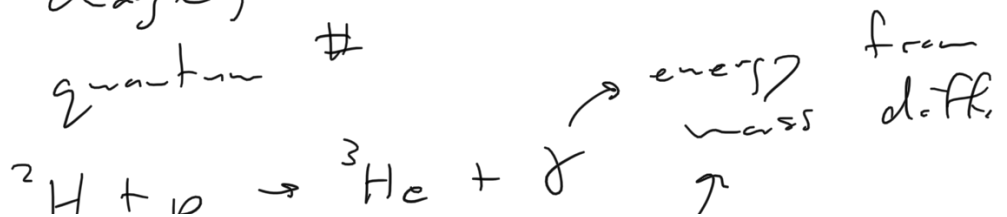
quantum tunneling

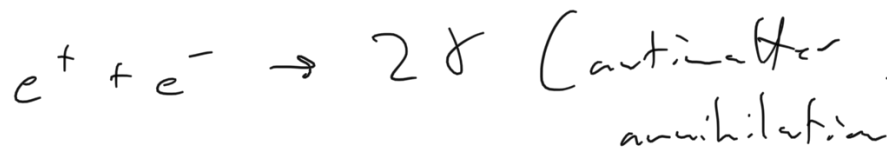
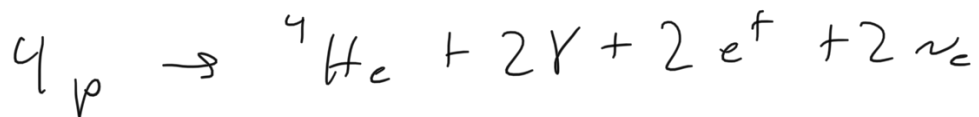
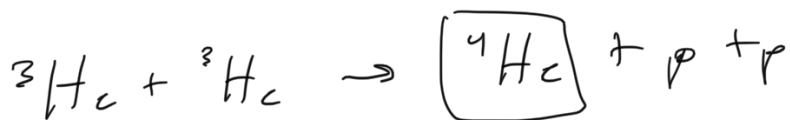
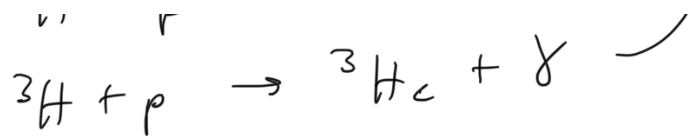
So, how exactly do these reactions create He in the core?

PP chain $T_c < 1.8 \times 10^7 \text{ K}$ (e.g., Sun)



$p \rightarrow n$, positron carries away extra charge, neutrino to conserve e^- quantum #





ν_e only interact via the weak force, which has a very ↓ cross-section, so Sun is essentially transparent to them & they escape

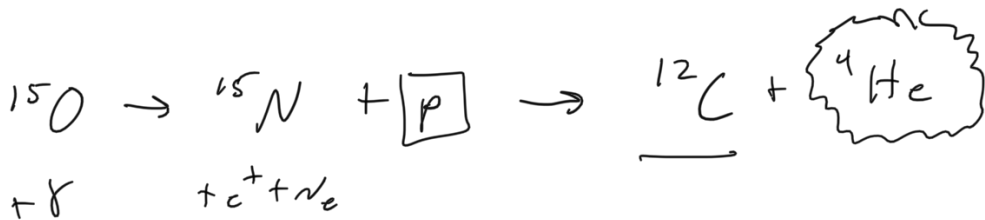
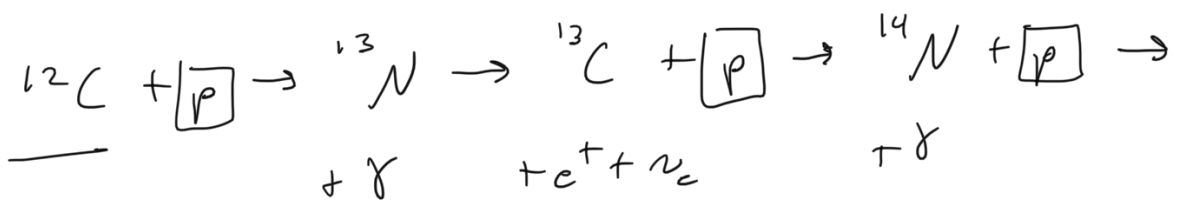
The γ -rays produced in the core take $\sim 10,000 - 100,000$ years to escape (b/c they scatter a lot - have a small m.f.p.). They are in LTE as they travel, so lose E as they travel, becoming visible IR at the surface.

powerous

$H \rightarrow He$ releases converts 0.7% of
the mass to energy
 $He \rightarrow$ bigger elements releases more!

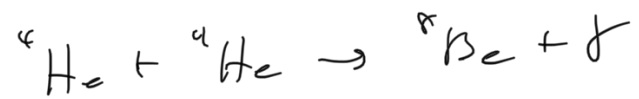
For more massive stars, where
 $T_c > 1.8 \times 10^8 K$, $H \rightarrow He$ via the
CNO cycle (C, N, O atoms used in
intermediary steps, but no net
change in their #s - they're
catalysts)

CNO cycle sets 1 additional γ -ray



To fuse $He \rightarrow C$, need the
triple alpha process

$T_c > 10^8 \text{ K}$ as well as dense environment,



↳ decays quickly,

but if it hits another ${}^4\text{He}$ first

