### Cosmic Microwave Background ASTR/PHYS 4080: Intro to Cosmology Week 8



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# **Brief History**

- 1934 (Richard Tolman) blackbody radiation in an expanding universe cools but retains its thermal distribution and remains a blackbody
- temperature of space of ~2.3K
- early universe (high expansion rate, assume matter domination)
- 1948 (Gamow) T~10<sup>9</sup>K when deuterium formed, argues for radiation domination in early universe; the existence of CMB
- 1948 (Alpher, Bethe, & Gamow [αβγ paper]), element synthesis in an expanding universe; calculations based on previous ideas
- dominated initial state); no mention of the observability. **ASTR/PHYS 4080: Introduction to Cosmology**

• 1941 (Andrew McKellar) excitation of interstellar CN doublet absorption lines gives an effective

• 1946 (Gamow) to match observed abundance, nuclei should be built up out of equilibrium in hot

1948 (Alpher & Herman) make corrections to previous results; state that present radiation temperature should be ~5K (close! but largely a coincidence; incorrect assumptions - neutron





# **Brief History**

- 1957 (Shmaonov) horn antenna at 3.2cm, find the absolute effective temperature of radio emission background 4±3K, independent of time and direction
- from helium abundance; realize Bell Labs telescope can constrain
- longer pure neutron initial state; weak interaction for neutron vs proton)
- 1965 (Dicke, Peebles, Roll, & Wilkinson) realize oscillating or singular universe might have thermal background; build detector to search; then they hear about its discovery
- 1965 (Penzias & Wilson) antenna has isotropic noise of 3.5±1.0K at wavelength of 7.35cm; white dielectric" generated by pigeons); explanation could be that of Dicke et al.

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• Early 1960s, (Zel'dovich, Doroshkevich, Novikov) estimate expected background temperature

1964 (Hoyle & Tayler) essentially correct version of primordial helium abundance calculation (no

careful experiment (e.g., shooed away pigeons roosted in the antenna; cleaned up "the usual





# **Brief History**

- with Penzias & Wilson for blackbody spectrum; isotropic to 10%
- 1966-1967 (Field & Hitchcock, Shklovsky, Thaddeus & Clauser, Thaddeus [following a CMB (McKellar's 1941 observation explained!)
- 1970s-1980s, ground, balloon, satellite observations
- isotropic blackbody and discovers the anisotropies.
- then the launches of WMAP (2001) and Planck (2009)

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• 1965 (Roll & Wilkinson) detect the radiation background at 3.2cm, with amplitude consistent

suggestion by Woolf]) independently show that the excitation of interstellar CN is caused by

• 1992, NASA's COsmic Background Explorer (COBE) satellite confirms CMB as nearly perfect

• Era of "precision cosmology" begins, especially with SNe measurements a few years later and





# Near perfect BB everywhere on the sky



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### Spatial variations on different scales



dT ~ 3.353 mK

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 $\frac{\delta T}{T}(\theta,\phi) \equiv \frac{T(\theta,\phi) - \langle T \rangle}{T}$ 

dT ~ 0.018 mK









# History of CMB space measurements



COBE 1990

WMAP 2003

Planck 2013

https://fineartamerica.com/featured/cosmic-microwave-background-radiation-carlos-clarivan.html

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### Primary aim to measure small-scale fluctuations



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# **Observing the CMB**



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# What produces the CMB and features we see?

- In the early universe, many interactions between particles (just like at the LHC) quarks, electrons, photons, neutrinos all transform into each other
- As universe expands, densities decrease and protons/electrons/photons dominate baryon soup
  - Eventually, electrons can be captured by protons to form atoms that are not immediately broken up by energetic photons -> recombination
  - Soon thereafter, the density of free electrons is too low to scatter photons, and the universe becomes transparent -> photon decoupling
- As the universe expands further, a time comes when a CMB photon scatters off an electron for one last time -> last scattering

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### Surface of Last Scattering



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### Recombination

### $\langle h\nu \rangle < H_{\rm ionize} = Q = 13.6 \text{ eV}$ $\langle h\nu \rangle = 2.7kT \text{ (BB spectrum)}$

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### $H + \gamma \rightleftharpoons p + e^{-}$

#### Implies $T \sim 60,000$ K —> much too high: BB spectrum has a tail

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### **Recombination** $H + \gamma \rightleftharpoons p + e^-$

$$n_{x}(p)dp = g_{x}\frac{4\pi}{h^{3}}\frac{p^{2}dp}{\exp([E-\mu_{x}]/kT)\pm 1} \qquad \text{(minus for bosons, plus for fermions)}$$

$$g \rightarrow 2 \text{ (for non-nucleons, g_{H}=4)}$$

$$chemical potential of photons = 0$$

$$\mu_{H} = \mu_{p} + \mu_{e}$$

$$n_{\gamma} = \frac{2.4041}{\pi^{2}} \left(\frac{kT}{\hbar c}\right)^{3} \qquad n_{x} = g_{x} \left(\frac{m_{x}kT}{2\pi\hbar^{2}}\right)^{3/2} \exp\left(\frac{-m_{x}c^{2} + \mu_{x}}{kT}\right)$$

$$\frac{n_{H}}{n_{p}n_{e}} = \frac{g_{H}}{g_{p}g_{e}} \left(\frac{m_{H}}{m_{p}m_{e}}\right)^{3/2} \left(\frac{kT}{2\pi\hbar^{2}}\right)^{-3/2} \exp\left(\frac{[m_{p} + m_{e} - m_{H}]c^{2}}{kT}\right) = \left(\frac{m_{e}kT}{2\pi\hbar^{2}}\right)^{-3/2} \exp\left(\frac{\mu_{e}kT}{4\pi\hbar^{2}}\right)^{-3/2} \exp\left($$

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Saha Equation





### **Recombination** $H + \gamma \rightleftharpoons p + e^{-1}$

$$\frac{n_H}{n_p n_e} = \frac{g_H}{g_p g_e} \left(\frac{m_H}{m_p m_e}\right)^{3/2} \left(\frac{kT}{2\pi\hbar^2}\right)^{-3/2} \exp\left(\frac{[m_p + m_e - m_H]c^2}{kT}\right) = \left(\frac{m_e kT}{2\pi\hbar^2}\right)^{-3/2} \exp\left(\frac{R}{kT}\right)$$
Defined as when protons and H atoms are equal:  

$$X \equiv \frac{n_p}{n_p + n_H} = 1/2$$

$$\frac{1 - X}{X} = n_p \left(\frac{m_e kT}{2\pi\hbar^2}\right)^{-3/2} \exp\left(\frac{Q}{kT}\right) \qquad \eta = \frac{n_p}{Xn_\gamma} \quad , \quad n_\gamma = \frac{2.4041}{\pi^2} \left(\frac{kT}{\hbar c}\right)^3$$
(set by current baryon/photon density)  

$$\frac{1 - X}{X^2} = 3.84\eta \left(\frac{kT}{m_e c^2}\right)^{3/2} \exp\left(\frac{Q}{kT}\right) = S \quad \longrightarrow \quad X = \frac{-1 + \sqrt{1 + 4S}}{2S} \quad \longrightarrow \quad kT_{\rm rec} = \frac{Q}{42}$$

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# Redshift of recomb., decoupling, & scattering



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recombination: z = 1380 when T = 3760K  $t_{age} = 250,000 \text{ yr}$ 

decoupling: when expansion rate surpasses scattering rate:  $\Gamma(z) = H(z)$  $1 + z = \frac{39.3}{X(z)^{2/3}}$ 

 $z \sim 1090$  (incl. non-eq. effects)

last scattering: when the optical depth is ~1 600  $\tau(t) = \int_t$  $\Gamma(t)dt$ 

redshift same as decoupling







### **Temperature fluctuations**



$$\ell = d_A \cdot \delta \theta = 12.8 \operatorname{Mpc}\left(\frac{\delta \theta}{1 \operatorname{rad}}\right) = 3.7 \operatorname{kpc}\left(\frac{\delta \theta}{1 \operatorname{arcmin}}\right)$$

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## **Spherical Harmonics & Power Spectrum**



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Represent function in terms of spherical harmonics

$$\frac{\delta T}{T}(\theta,\phi) = \sum_{\ell=0}^{\infty} \sum_{m=\ell}^{\ell} a_{\ell m} Y_{\ell m}(\theta,\phi)$$
sum Y over  
m, get  
Legend  
polyne  

$$C(\theta) = \left\langle \frac{\delta T}{T}(\hat{n}) \frac{\delta T}{T}(\hat{n}') \right\rangle_{\hat{n} \cdot \hat{n}' = \cos \theta}$$

$$= \frac{1}{4\pi} \sum_{\ell=0}^{\infty} (2\ell + 1) C_{\ell} P_{\ell}(\cos \theta)$$





# (2 point) Correlation function



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## Where do the peaks come from?

First, let's define the scale at which pieces of the universe could be in causal contact

At last scattering, universe evolves as if there's only radiation and matter, so we can easily calculate the horizon distance

$$d_{\rm hor}(t_{\rm ls}) = a(t_{\rm ls})c \int_0^{t_{\rm ls}} \frac{dt}{a(t)} = 2.24ct_{\rm ls} \approx 250 \ \rm kpc$$

By definition, the angular scale this occurs at is given by the angular diameter distance

$$\theta_{\rm hor} = \frac{d_{\rm hor}(t_{\rm ls})}{d_A} = \frac{250 \text{ kpc}}{12.8 \text{ Mpc}} \approx 1.1^{\circ}$$

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# Sachs-Wolfe Effect



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## Acoustic peaks



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# CMB provides a giant triangle of known size!



a If universe is closed, "hot spots" appear larger than actual size





b If universe is flat,
 "hot spots" appear
 actual size

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c If universe is open, "hot spots" appear smaller than actual size



## Acoustic peaks



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# Acoustic peaks



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### https://lambda.gsfc.nasa.gov/ education/cmb\_plotter/

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# Fitting the power spectrum in detail yields narrow constraints

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