

ASTR 5590 - Black Holes

★ Refer to review sheets on BHs

Classically defined as having an escape velocity $>$ speed of light

$$E_{esc} = \frac{1}{2} m v_{esc}^2 - \frac{GMm}{r} = 0$$

$$v_{esc} = \sqrt{\frac{2GM}{r}} = c \quad \left(\text{for a BH,} \right)$$

$$\text{so } r_s = \frac{2GM}{c^2}$$

→ light can't escape if inside this radius, so object smaller than this size must be black

Just happens to be radius of such an object in GR, derived from the metric: $R_s = \frac{2GM}{c^2} = 3 \text{ km} \left(\frac{M}{M_\odot} \right)$

Schwarzschild Radius

→ for a point mass

— — — — —

In GR, V_S applies to radius r —
 spacetime curvature prevents particles
 from moving radially away

$$ds^2 = \alpha dt^2 - \frac{1}{c^2} \left[\alpha^{-1} dr^2 + r^2 (d\theta^2 + \sin^2\theta d\phi^2) \right]$$

$$\alpha = 1 - \frac{2GM}{rc^2}$$

$ds \rightarrow$ distance in space-time b/w
 2 events (invariant regardless
 of definition of coords)

Special Relativity, use the Minkowski
 metric

$$ds^2 = dt^2 - \frac{1}{c^2} \left[dr^2 + r^2 (d\theta^2 + \sin^2\theta d\phi^2) \right]$$

Far from the point mass, $r \rightarrow \infty$,
 $\alpha \rightarrow 1$ & metrics are the same


Nearer the mass, the metrics give

$$= dt' \frac{1}{\alpha^{1/2}} \left[\alpha - \frac{1}{c^2 \alpha} \frac{v^2}{2} dt'^2 \right]$$

$$= dt' \left[1 - \frac{dr^2}{c^2 \alpha} \frac{1}{dt'^2} \right]^{1/2}$$

An increment of "geodesic distance" is actually $dx = \alpha^{-1/2} dr$

$$d\tau = dt' \left[1 - \frac{1}{c^2} \frac{dx^2}{dt'^2} \right]^{1/2} = dt_{\text{as}}$$



 Lorentz factor γ

 @ v is moving thru S

 (as frame)

↳ time dilation

$$\text{If } r \leq r_g = \frac{2GM}{c^2}, \quad v_{\infty} \rightarrow 0$$

★ no information can be transferred
 - also no light

GR also tells us that KE of a particle contributes to the grav force felt, so a particle w/ any mass close to a BH will eventually reach distance where this "extra gravity"

a ...
becomes larger than the centrifugal force, so it must fall in it

$$l = m v r \leq 2 r_g c m$$

Last stable circular orbit @ $r = 3 r_g$

There is NO singularity @ r_g : just looks that way since choice of coords in metric

There IS a physical singularity @ $r = 0$

↳ likely a conseq. of GR not being quantized

HERE

Real BHs are almost certainly rotating described by the Kerr metric

Parameterize ang. momentum J by

$$a = \frac{J}{Mc} \quad (\text{units of distance})$$

event horizon becomes

$$r_{\pm} = \frac{GM}{c^2} \pm \left[\left(\frac{GM}{c^2} \right)^2 - a^2 \right]^{1/2}$$

Minimal spinning BH has $a = \frac{GM}{c^2}$

(maximum))
[Observationally, "a" is in units of $\frac{GM}{c^2}$
so its maximum value is 1]
→ half the size of a non-spinning BH

"Frame dragging" → particles are forced
to move in direction of spin

B/c event horizon smaller, particles can
orbit closer, thus more of their
rest mass E can be radiated during inf.

→ non-rotating: $\sim 6\%$ of mc^2
(binding E of particle @ $r = 3r_g$)

→ for particles orbiting in same direction
as rotation, can get $\sim 40\%$ of mc^2
for $a = 1$ spinning BH

Inside the "static" radius, where particles
cannot have $d\theta = 0$, is called ergosphere

- if a particle ends up on a
'negative energy' orbit inside
then splits into 2 w/ one falling in
& the other escaping, the latter

particle can leave w/ more E
than the original particle fell in;
↳ garbage dump power for an advanced
alien civ. → gain up to 30% of
 mc^2
↳ called the Penrose process

$$\text{Hawking radiation: } T = \frac{\hbar c^3}{8\pi G M k} \approx \frac{10^{-7}}{(M/M_\odot)} \text{ K}$$

→ merging QM near event horizon