

# ASTR 5590 - Chapter 11d

## Ideal Gas Review $p = (\gamma - 1) \epsilon$

A perfect gas relates  $p$  &  $\epsilon$  ( $\epsilon$  dens.)  
by the adiabatic index  $\gamma$  (ratio of heat capacities)

What is  $\gamma$ ?

ratio of heat capacities  
@ const.  $p$  & const.  $V$   $\rightarrow \epsilon$  to raise unit mass  $1 \Delta T$  of temp

$$\gamma = \frac{c_p}{c_v}$$

Ideal Gas  $\rightarrow pV^\gamma = \text{const.}$  : defines an adiabatic

$\rightarrow p \propto \rho^\gamma$  (no heat added to system)

Like before w/ MHD waves, can have small perturbations  $p = p_0 + \delta p$ ,  $\frac{\delta p}{p} \ll 1$ , expand to 1st order, get wave eqns

Sound speed is  $c_s = \sqrt{\left(\frac{\partial p}{\partial \rho}\right)_s}$  (Newton's law plus conservation mass)

or since  $\frac{\partial p}{\partial \rho} = \gamma p^{\gamma-1} = \gamma \frac{p}{\rho}$ ,  $c_s = \sqrt{\gamma \frac{p}{\rho}}$

|          | Ideal Gas<br>non-rel. gas | Radiation<br>relativistic gas |
|----------|---------------------------|-------------------------------|
| $\gamma$ | 5/3                       | 4/3                           |
| $p$      | $\rho kT$                 | $\frac{4}{3} \sigma T^4$      |

|     |                                       |                            |
|-----|---------------------------------------|----------------------------|
|     | $\mu m_p$                             | $\frac{3}{2} \frac{kT}{c}$ |
| $E$ | $\frac{3}{2} \frac{\rho kT}{\mu m_p}$ | $\frac{4\sigma}{c} T^4$    |

$\sigma \rightarrow$  Stefan-Boltzmann const.

$$= 5.67 \times 10^{-5} \text{ erg / (cm}^2 \text{K}^4 \text{s)}$$

$\mu m_p \rightarrow$  mean mass per particle

$\mu = 1$  for neutral H

$= \frac{1}{2}$  for 100% ionized H

Here  $\rightarrow$

## Shock Waves

Figs. 11.6

In frame of shock (b), situation is

in steady state:  $\frac{\partial}{\partial t} = 0$

Either side of shock, quantities conserved:

mass | energy | momentum

Mass) Only change in  $x$  direction ( $y, z$  <sup>unaffected</sup>)

flux of particles thru an area

# dens. = flux  $\cdot \frac{1}{v}$ , @ shock



flux in = flux out, so  $\boxed{\rho_1 v_1 = \rho_2 v_2}$

Energy fluid dynamics: E flux thru  
a surface normal to  $\vec{v}$  is  $\rho \vec{v} (\frac{1}{2} v^2 + w)$

$w \equiv$  enthalpy per unit mass

$$= E_{in} + PV$$

↑ internal energy per unit mass  
← vol. per unit mass

conservation gives

$$\rho_1 v_1 (\frac{1}{2} v_1^2 + w_1) = \rho_2 v_2 (\frac{1}{2} v_2^2 + w_2)$$

$$v_1 (\frac{1}{2} \rho_1 v_1^2 + E_1 + P_1) = v_2 (\frac{1}{2} \rho_2 v_2^2 + E_2 + P_2)$$

Momentum  $P_1 + \rho_1 v_1^2 = P_2 + \rho_2 v_2^2$

- regular momentum  $mv$ , plus pressure  
(force per area)

↳ think of  $P$  as pushing gas

These 3 called Shock Jump

Conditions

Assuming adiabatic conditions on  
both sides (but different adiabats)

rewrite energy conservation using

$$E/v = \left( \frac{1}{\gamma-1} \right) p$$

$$\boxed{v_1^2 + \frac{2\gamma}{\gamma-1} \frac{p_1}{\rho_1} = v_2^2 + \frac{2\gamma}{\gamma-1} \frac{p_2}{\rho_2}}$$

Can define Mach #  $\Rightarrow$   $\frac{\text{velocity of shock}}{\text{sound speed in pre-shock gas}}$

$$M_1 = \frac{v_s}{c_{s,1}}, \quad M^2 = \frac{\rho_1}{\gamma p_1} v_s^2, \quad v_1 = |v_s|$$

From mass cons.  
& E cons.

$$\boxed{\frac{v_2}{v_1} = \frac{\rho_1}{\rho_2} = \frac{2}{\gamma+1} \frac{1}{M_1^2} + \frac{\gamma-1}{\gamma+1}}$$

From mom. cons.  
& E cons.

$$\boxed{\frac{p_2}{p_1} = \frac{2\gamma M_1^2 - (\gamma-1)}{\gamma+1}}$$

Using Ideal Gas  
 $p = \rho kT$

$$\boxed{\frac{T_2}{T_1} = \frac{[2\gamma M_1^2 - (\gamma-1)][2 + (\gamma-1)M_1^2]}{(\gamma+1)^2 M_1^2}}$$

$$M_1^2 \geq 1$$

Consider Strong Shocks,  $M \gg 1$ , above  
reduces to

$$\underline{\gamma = 5/3}$$

$$\frac{r_2}{r_1} = \frac{2r}{\gamma+1} M^2 \longrightarrow \frac{5}{4} M^2$$

$$\frac{\rho_2}{\rho_1} = \frac{\gamma+1}{\gamma-1} \longrightarrow \underline{4} \star$$

$$\frac{T_2}{T_1} = \frac{2\gamma(\gamma-1)M_1^2}{(\gamma+1)^2} \longrightarrow \frac{5}{16} M_1^2$$