

ASTR 5590 - Chapter 9c

Comptonization \rightarrow when Compton scattering dominates evolution of spectrum

Assume non-rel., $1/2 kT_e \ll mc^2$ &
 $E = h\nu = \hbar\omega \ll mc^2$
 \rightarrow Thomson scat. valid

$$\frac{\Delta E}{E} = \frac{h\nu}{mc^2} (1 - \cos\theta), \text{ avg. over } \theta$$

$$\left\langle \frac{\Delta E}{E} \right\rangle_{e^-} = \frac{h\nu}{mc^2} \rightarrow E \uparrow \text{ of } e^-$$

$$IC \rightarrow \frac{dE}{dt} = \frac{4}{3} \sigma_T c u_{\text{rad}} \left(\frac{v^2}{c^2} \right) \quad \begin{matrix} \nearrow 1 \\ \swarrow \text{cross-section} \\ \searrow \text{speed} \end{matrix} \quad \begin{matrix} \text{(prod. of collision)} \\ \text{density} \end{matrix}$$

$$N_{\text{ph}} = \# \text{ photons scattered per sec} \rightarrow n_{\text{ph}} \sigma_T c$$

$$n_{\text{ph}} = \frac{u_{\text{rad}}}{h\nu} \quad \leftarrow \# \text{ density}$$

$$\frac{dE}{dt} = \frac{4}{3} N_{\text{ph}} h\nu \left(\frac{v^2}{c^2} \right)$$

So, ΔE gained per photon is, on average

$$\left\langle \frac{\Delta E}{E} \right\rangle = \frac{4}{3} \frac{v^2}{c^2}$$

If e^- are in thermal eq., $\frac{1}{2} m_e \langle v^2 \rangle = \frac{3}{2} kT_e$

$$\text{so } \left\langle \frac{\Delta E}{E} \right\rangle_{\nu} = \frac{4kT_e}{m_e c^2} \rightarrow E \uparrow \text{ photon}$$

$$\begin{aligned} \text{Total } \left\langle \frac{\Delta E}{E} \right\rangle_{\nu, \text{net}} &= \left\langle \frac{\Delta E}{E} \right\rangle_{\nu} - \left\langle \frac{\Delta E}{E} \right\rangle_{e^-} \\ &= \frac{4kT_e - h\nu}{m_e c^2} \end{aligned}$$

If $h\nu > 4kT_e \rightarrow E$ transferred to e^-
 $h\nu < 4kT_e \rightarrow E$ transferred to γ

Basically, more energetic comp. gives some E to less energetic one

Consider case where e^- hotter, $4kT_e \gg h\nu$

$$\left\langle \frac{\Delta E}{E} \right\rangle_{\text{per collision}} = \frac{4kT_e}{m_e c^2}$$

- a region ν /size l + # dens. n_e has
 an optical depth $\tau_e = n_e \sigma_T l$

If $\tau_e \gg 1$, random walk for photons escaping
 $l \sim N^{1/2} \lambda_e$ ($\lambda_e = [n_e \sigma_T]^{-1}$)
 mfp \uparrow

$$N = \left(\frac{l}{\lambda_e} \right)^2 = \left(\frac{\tau_e}{n_e \sigma_T} n_e \sigma_T \right)^2 = \tau_e^2$$

If $\tau_e \ll 1$, basically only get 1 scat.
 or less, $\nu/N = \tau_e$

For a spectrum to be distorted by Comptonization, $4\gamma \gtrsim 1$ where

$$\gamma = \frac{kT_e}{m_e c^2} \max(\bar{v}_e, \bar{v}_e^2)$$

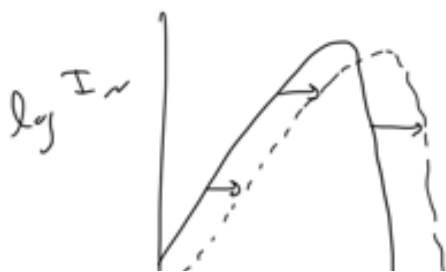
$\uparrow < 1$ $\uparrow > 1$

γ parameter: Compton optical depth

If $\gamma \gg 1$, photons achieve BB spectrum

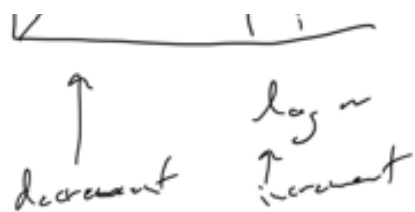
If $\gamma \ll 1$, then $\gamma = \frac{kT_e}{m_e c^2} n_e \sigma_T l$

- consider e^- hotter than radiation
(cluster gas + CMB photons)



Produce a distortion of the spectrum

Ch F: 9.11



$$\frac{\Delta I_r}{I_r} = -2\gamma$$

Clusters of galaxies, $n/T_e \sim 10^8 \text{K}$,
 $\gamma \approx 10^{-5}$, so change in CMB small

If integrate over entire cluster, defining

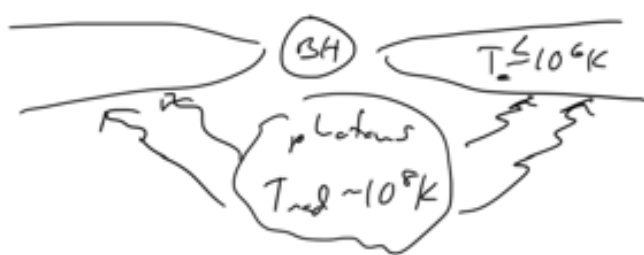
$$\gamma = \frac{kT_e}{m_e c^2} \sigma_T \int n_e dl$$

$$\rightarrow Y = \int \gamma dA = \frac{kT_e}{m_e c^2} \sigma_T \int n_e dl \cdot dA$$

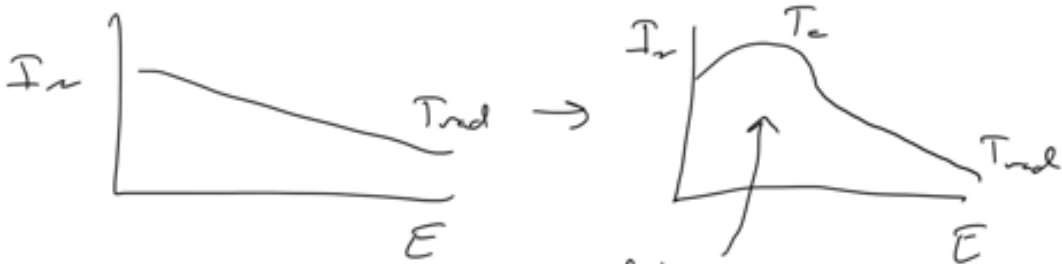
$$\rho = n_e kT$$

$$= \frac{\sigma_T}{m_e c^2} \int n_e kT_e dV \approx \frac{\sigma_T}{m_e c^2} PV \quad \begin{matrix} \text{thermal} \\ E \text{ of} \\ \text{cluster} \\ \text{gas} \end{matrix}$$

What if photons have TE than e^- ?



- Give E to e^-
- Create a "Compton Heap"



soft X-rays are Comptonized

